

OPEN ACCESS

Application of Fibonacci oscillators in the Debye model

To cite this article: André A A Marinho *et al* 2014 *J. Phys.: Conf. Ser.* **568** 012009

View the [article online](#) for updates and enhancements.

Related content

- [NMR study on kagome-lattice antiferromagnet \$\text{Cs}_2\text{Mn}_3\text{LiF}_6\$](#)
S Moriya, Y Usuf, H Tanaka *et al.*
- [Non-extensive entropy of bosonic Fibonacci oscillators](#)
Abdullah Algin
- [Influence of thermal gradient in vortex states of mesoscopic superconductors](#)
E C S Duarte, A Presotto, D Okimoto *et al.*

Recent citations

- [Thermostatistics of bosonic and fermionic Fibonacci oscillators](#)
Abdullah Algin *et al*
- [Effective approach for taking into account interactions of quasiparticles from the low-temperature behavior of a deformed fermion-gas model](#)
Abdullah Algin and Ali Serdar Arıkan
- [Thermal and electrical properties of a solid through Fibonacci oscillators](#)
André A A Marinho *et al*

Application of Fibonacci oscillators in the Debye model

André A.A. Marinho^{a,b}, F.A. Brito^b, C. Chesman^a

^a Departamento de Física Teórica e Experimental, Universidade Federal do Rio Grande do Norte, Caixa Postal 1634, 59072-970 Natal, Rio Grande do Norte, Brazil

^b Departamento de Física, Universidade Federal de Campina Grande, 58109-970 Campina Grande, Paraiba, Brazil

E-mail: andre@dfte.ufrn.br or andreofisico@hotmail.com.br

Abstract. In this paper we study the thermodynamics of a crystalline solid by applying q -deformed algebra of Fibonacci oscillators through the generalized Fibonacci sequence of two real and independent deformation parameters q_1 and q_2 . We find a (q_1, q_2) -deformed Hamiltonian and consequently the q -deformed thermodynamic quantities. The results led us to interpret the deformation parameters acting as disturbance or impurities factors modifying the characteristics of a crystal structure. More specifically, we found the possibility of adjusting the Fibonacci oscillators to describe the change of thermal conductivity of a given element as one inserts impurities.

1. Introduction

The study of quantum groups and quantum algebras has attracted great interest in recent years and stimulated intense research in various fields of physics [1, 2, 3] taking into account a range of applications covering astrophysics and condensed matter, for instance, black holes and high-temperature superconductors [4].

A possible mechanism to generate a deformed version of the classical statistical mechanics consists in replacing the Gibbs-Boltzmann distribution by a deformed version. In this sense it is postulated a form of deformed entropy which implies a generalized theory of thermodynamics [5, 6, 7, 8]. In this context, it was demonstrated that a natural realization of the q -deformed thermodynamics can be done via q -calculation formalism [9, 10].

A new proposal for the q -calculation is the inclusion of two distinct deformation parameters in some physical applications. Starting from the generalization of q -algebra [11], in Ref. [12] was generalized the Fibonacci sequence. Here, the numbers are in that sequence of generalized Fibonacci oscillators, where new parameters (q_1, q_2) are introduced [13, 14, 15]. They provide a unification of quantum oscillators with quantum groups, keeping the degeneration property of the spectrum invariant under the symmetries of the quantum group. The quantum algebra with two deformation parameters may have a greater flexibility when it comes to applications in realistic phenomenological physical models [16, 17].

We know that a solid is formed by a large number of atoms bound by cohesion forces of various types. The motion of atoms in a solid is very narrow, causing each atom to move only within a small neighborhood, executing vibratory motion around its equilibrium point. In a crystalline solid, the equilibrium points of atomic vibrations form a regular spatial structure, a cubic structure, for example.

The study conducted by Anderson, Lee and Elliot [18, 19, 20] shows that the presence of defects or impurities in a crystal modifies the electrostatic potential in their neighborhood, breaking the translational



symmetry of the periodic potential. This perturbation can produce electronic wave functions located near the impurity, ceasing to be propagated throughout the crystal.

The conductivity of semiconductors can be dramatically altered by the presence of impurities, i.e., different from atoms that make up the pure crystal. This property allows the production of a variety of electronic devices of the same semiconductor material. This process of placing impurities in semiconductor materials is called doping.

In the present study we follow the lines of [21, 22, 23] to address the issues of (q_1, q_2) -deformed thermodynamics through application of a generalized q -algebra to a Debye solid. We find that the deformation is related to phenomena due to impurities or disorder factors in the system. We shall mainly show that the Fibonacci oscillators may act as defects or impurities in the crystal lattice, allowing us to modify quantities such as Debye temperature, thermal and electrical conductivities.

2. Algebra of the Fibonacci oscillators

The generalization of integers usually is given by a sequence. The two well-known ways to describe a sequence are the arithmetic and geometric progressions. However, the Fibonacci sequence encompasses both. Generalizing this sequence, we get the Fibonacci oscillators, so the spectrum can now be given by the Fibonacci integer.

Defining the Fibonacci *basic number*, the q -deformed quantum oscillator is now defined by the Heisenberg algebra in terms of the annihilation and creation operators c and c^\dagger , respectively, and the number operator N , as follows:

$$[x_{i,q_1,q_2}] = c_i^\dagger c_i = \frac{q_2^{2x_i} - q_1^{2x_i}}{q_2^2 - q_1^2}, \quad (1)$$

$$c_i c_i^\dagger - q_1^2 c_i^\dagger c_i = q_2^{2N_i}, \quad \text{and} \quad c_i c_i^\dagger - q_2^2 c_i^\dagger c_i = q_1^{2N_i}, \quad (2)$$

$$[1 + x_{i,q_1,q_2}] = q_1^2 [x_{i,q_1,q_2}] + q_2^{2N_i}, \quad \text{or} \quad [1 + x_{i,q_1,q_2}] = q_2^2 [x_{i,q_1,q_2}] + q_1^{2N_i}. \quad (3)$$

Let us now generalize to (q_1, q_2) -deformed algebra the study we accomplished in [22, 23] for only one q -deformation parameter. Now we have the following Hamiltonian with the respective energy eigenvalues

$$\mathcal{H}_{q_1,q_2} = \frac{\hbar\omega}{2} (cc^\dagger + c^\dagger c), \quad (4)$$

$$E_{n_{i,q_1,q_2}} = \frac{\hbar\omega_E}{2} ([n_{i,q_1,q_2}] + [n_{i,q_1,q_2} + 1]) = \frac{\hbar\omega_E}{2} + \frac{\hbar\omega_E (2 \ln(q_2) - 2 \ln(q_1)) n}{q_2^2 - q_1^2}. \quad (5)$$

Notice that when $q_1 = 1$ and $q_2 \rightarrow 1$ (and vice-versa) the deformation vanishes, so that we have the usual definition

$$E_n = \frac{\hbar\omega_E}{2} (2n + 1). \quad (6)$$

3. (q_1, q_2) -deformed Debye Solid

Corrections of Einstein's model are given by the Debye model, allowing us to integrate from a continuous spectrum of frequencies up to the Debye frequency ω_D , as the number of normal modes of vibration is $3N$, as is found in the literature [24, 25, 26]. Following the same mathematical development, and inserting the given parameters (q_1, q_2) , we can write the specific heat for any temperature in the form:

$$c_{V_{q_1,q_2}}(T) = 3\kappa_B D(\alpha_{0_{q_1,q_2}}), \quad (7)$$

where $D(\alpha_{0_{q_1,q_2}})$ is the (q_1, q_2) -deformed Debye function, defined by

$$D(\alpha_{0_{q_1,q_2}}) = \frac{3}{(\alpha_{0_{q_1,q_2}})^3} \int_0^{\alpha_{0_{q_1,q_2}}} \frac{\alpha^4 \exp(\alpha)}{[\exp(\alpha) - 1]^2} d\alpha, \quad (8)$$

$$\alpha_{0_{q_1, q_2}} = \frac{\hbar \omega_{D_{q_1, q_2}}}{\kappa_B T} = \frac{\theta_{D_{q_1, q_2}}}{T}, \quad \omega_{D_{q_1, q_2}} = 2\omega_D \left(\frac{\ln(q_2) - \ln(q_1)}{q_2^2 - q_1^2} \right). \quad (9)$$

where $\omega_{D_{q_1, q_2}}$ and $\theta_{D_{q_1, q_2}}$, are the (q_1, q_2) -deformed Debye frequency and temperature, respectively, and ω_D is the Debye frequency characteristic. Integrating Eq.(8) by parts one finds

$$D(\alpha_{0_{q_1, q_2}}) = -\frac{3\alpha_{0_{q_1, q_2}}}{\exp(\alpha_{0_{q_1, q_2}}) - 1} + \frac{12}{\alpha_{0_{q_1, q_2}}^3} \left\{ -\frac{\pi^4}{15} - \frac{\alpha_{0_{q_1, q_2}}^4}{4} + \alpha_{0_{q_1, q_2}}^3 \ln[1 - \exp(\alpha_{0_{q_1, q_2}})] \right. \\ \left. + 3\alpha_{0_{q_1, q_2}}^2 Li_2[\exp(\alpha_{0_{q_1, q_2}})] - 6\alpha_{0_{q_1, q_2}} Li_3[\exp(\alpha_{0_{q_1, q_2}})] + 6Li_4[\exp(\alpha_{0_{q_1, q_2}})] \right\}, \quad (10)$$

where

$$Li_n(z) = \sum_{k=0}^{\infty} \frac{z^k}{k^n}, \quad (11)$$

is the polylogarithm function. We can check that, as in the usual Debye solid, the low-temperature specific heat in a q -deformed Debye solid is proportional to T^3 , due to phonon excitation, a fact that is in agreement with experiments. Thus, let us express the (q_1, q_2) -deformed specific heat for low temperatures as follows:

$$c_{V_{q_1, q_2}} = \frac{12\pi^4 \kappa_B}{5} \left(\frac{T}{\theta_{D_{q_1, q_2}}} \right)^3 = 1944 \left(\frac{T}{\theta_{D_{q_1, q_2}}} \right)^3 \frac{J}{molK}. \quad (12)$$

Through the settings for the thermal conductivity (κ) and electrical (σ) — see [27], we can rewrite the (q_1, q_2) -deformed specific heat in the useful form

$$\kappa_{q_1, q_2} = \frac{\kappa c_{V_{q_1, q_2}}}{c_V} \quad \text{and} \quad \sigma_{q_1, q_2} = \frac{\kappa_{q_1, q_2} \sigma}{\kappa}. \quad (13)$$

For illustration purposes, we choose iron (Fe) and chromium (Cr), two materials that can be employed in many areas of interest. In Figs.(1) we present deformed values of Fe (Fe_{q-def}) (black curve) and Cr (Cr_{q-def}) (green curve), for values $q_1 = 1$ and $q_2 = 0.1, \dots, 1$, where for this range we assume the maximum deformation ($q_1 = 1$ and $q_2 = 0.1$) and the pure element (bulk) ($q_1 = 1$ and $q_2 = 1$). The other elements are represented by colors and indicated in the very figure.

On the left side of the Fig.(1), we can observe that before reaching their limits, black and green curves can assume the values of Debye temperatures (θ_D) of other elements. The Fe_{q-def} e.g., equates to: beryllium (Be) when $q_2 \approx 0.23$, chromium (bulk) (Cr) $q_2 \approx 0.75$ and osmium (Os) $q_2 \approx 0.94$. On the right, we have the behavior of the curves obtained for the specific heat c_V . We note that the behavior is quite different from the previous curves θ_D , i.e., the curves start at lower values (maximum deformation) until they reach their pure values. Having Fe_{q-def} as an example again, it is possible to see, as it reaches the value of specific heat capacity of all the elements, including Cr (bulk) when $q_2 \approx 0.74$.

Let us now return to Eq.(10), where we have the complete Debye function $D(\alpha_{0_{q_1, q_2}})$ to see the behavior of the thermal conductivity. Thus, in the Fig.(2), we show a comparison to the thermal conductivity κ for a pure and impure material. We have the thermal conductivity as a function of temperature (T) for Fe (bulk) and a combination of the values q_1 and q_2 for the Fe_{q-def} (impure). Therefore, we have the Fe (red) when $q_1 = q_2 = 1$ (pure), and when $q_1 = 0.2$ and $q_2 = 1.8$, we have a similar behavior to silicon (Si) (golden), whose value for the Debye temperature is $\theta_D = 645K$. Finally, for a combination of values $q_1 = 1$ and $q_2 = 1.4$ we have a curve that is similar to that of zinc (Zn) (orange), with $\theta_D = 327K$.

One should note that deformation is clearly playing the role of impurity concentration in the material sample. This is because deformation acts directly on the Debye temperature, which means that the Debye frequency is modified. Changing the Debye frequency is a clear sign of the material being modified by impurities.

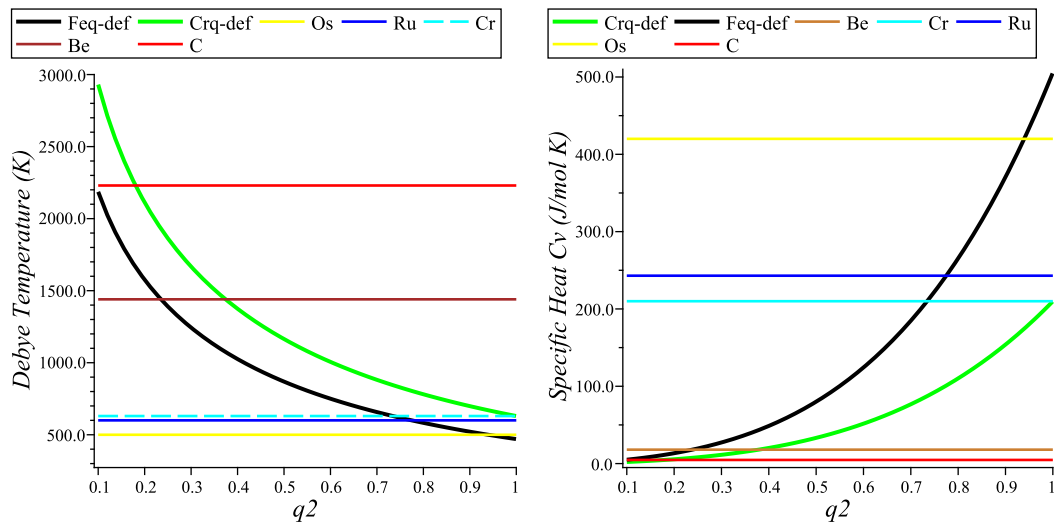


Figure 1. Debye temperature θ_D depending on the variation $q_2 = 0.1, \dots, 1$ and $q_1 = 1$, (**left**). Specific Heat c_V depending on the variation $q_2 = 0.1, \dots, 1$ and $q_1 = 1$, (**right**)

4. Conclusions

The initial idea that q -algebra acts as a factor of disorder or impurity is enhanced by inserting two factors, the so-called Fibonacci oscillators. In Figs.(1,2), we note that the elements that suffer deformation may become similar to others. The existence of more degrees of freedom as in the present case of two deformation parameters, q_1 and q_2 , can be well associated with different types of deformations related to two distinct phenomena of disorders or impurities such as, for instance, one due to pressure generating disorders and other due to doping, respectively.

Acknowledgments

We would like to thank CNPq, CAPES, and PNPd/PROCAD-CAPES, for partial financial support.

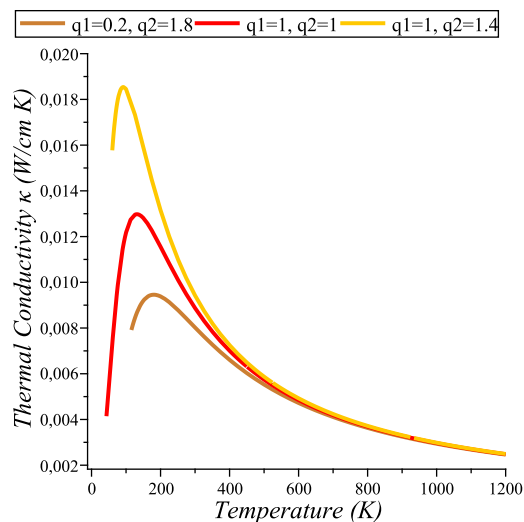


Figure 2. Thermal conductivity κ of Fe as a function of temperature in the range of $T = 0, \dots, 1200K$, and values for q_1 and q_2 given explicitly.

References

- [1] L. Biedenharn, J. Phys. A: Math. Gen. **22**, L873 (1989).
- [2] A. Macfarlane, J. Phys. A: Math. Gen. **22**, 4581 (1989).
- [3] A.U. Klimyk, *Spectra of Observables in the q-Oscillator and q-Analogue of the Fourier Transform*, Methods and Applications, **1**, 8, (2005).
- [4] F. Wilczek, *Fractional Statistics and Anyon Superconductivity*, World Scientific, Singapore, (1990).
- [5] C. Tsallis, J. Stat. Phys. **52**, 479 (1988).
- [6] G. Kaniadakis, M. Lissia, A. M. Scarfone, Phys. Rev. E **71**, 046128 (2005).
- [7] S. Abe, Phys. Lett. A **224**, 326 (1997).
- [8] A. Lavagno, Phys. Lett. A **301**, 13-18 (2002).
- [9] A. Lavagno, A. M. Scarfone e P. N. Swamy, J. Phys. A: Math. Theor. **40**, 8635-8654 (2007).
- [10] Liu Hui, et al., Whuan Univ. J. Nat. Sciences **15**, 57-63 (2010).
- [11] F.H. Jackson, Proc. Edin. Math. Soc. **22**, 28-39(1904).
- [12] M. Arik, et al., Z. Phys. C **55**, 89-95 (1992).
- [13] A. Algin, Phys. Lett. A **292**, 251-255 (2002);
A. Algin, B. Deviren, J. Phys. A: Math. Gen. **38**, 5945-5956 (2005).
A. Algin, J. Stat. Mech. Theor. Exp. **P10009**, 10 (2008).
A. Algin, E. Arslan, J. Phys. A: Math. Theor. **41**, 365006 (2008).
A. Algin, E. Arslan, Phys. Lett. A **372**, 2767-2773 (2008).
Algin A, Arik M., Kocabicakoglu D., Int. J. Theor. Phys.**47**, 1322-1332 (2008).
A. Algin, J. Stat. Mech. Theor. Exp. **P04007**, 04 (2009).
A. Algin, J. CNSNS **15**, 1372-1377 (2010).
- [14] A.M. Gavrilik, A.P. Rebish, Mod. Phys. Lett. A **22**, 949-960 (2007).
- [15] A.A. Marinho, F.A. Brito, C. Chesman, Physica A **411**, 74-79 (2014).
- [16] Daoud M., Kibler M., Phys. Lett A **206**, 13-17 (1995).
- [17] Gong R S, Phys. Lett A **199**, 81-85 (1995).
- [18] P.W. Anderson, Phys. Rev. **5**, 109 (1958).
- [19] P.A. Lee, T.V. Ramakrishnan, Rev. Mod. Phys. **2**, 57 (1985).
- [20] Elliott et al., Rev. Mod. Phys. **3**, 46 (1974).
- [21] F.A. Brito, A.A. Marinho, Physica A **390**, 2497-2503 (2011). .
- [22] A.A. Marinho, F.A. Brito, C. Chesman, Physica A **391**, 3424-3434 (2012).
- [23] D. Tristant, F.A. Brito, Physica A **407**, 276-286 (2014).
- [24] R.K. Pathria, *Statistical Mechanics*, Pergamon press, Oxford (1972)
- [25] K. Huang, *Statistical Mechanics*, John Wiley & Sons, (1987)
- [26] C. Kittel, *Introduction to Solid State Physics*, John Wiley & Sons, (1996)
- [27] J.M. Ziman, *Electron and Phonons - The Theory of Transport Phenomena in Solids*, Oxford Univ. Press, (1960).